



THE PHYSICS OF FOAM

- **Boulder School for Condensed Matter and Materials Physics**
July 1-26, 2002: Physics of Soft Condensed Matter

1. Introduction

Formation

Microscopics

2. Structure

Experiment

Simulation

3. Stability

Coarsening

Drainage

4. Rheology

Linear response

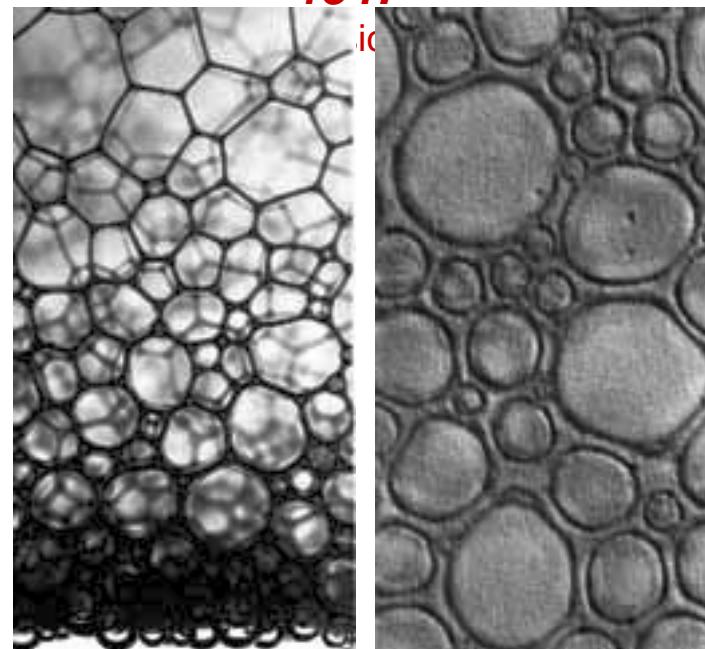
Rearrangement & flow

Douglas J. DURIAN

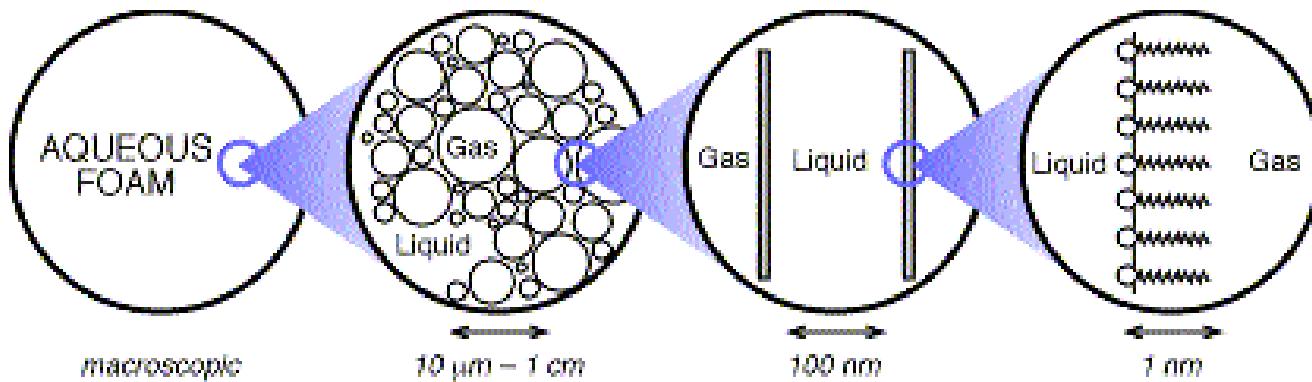
UCLA Physics &

Astronomy

*Los Angeles, CA 90095-
1547*



- ...a random packing of bubbles in a relatively small amount of liquid containing surface-active impurities
 - Four levels of structure:



- Three means of time evolution:
 - Gravitational drainage
 - Film rupture
 - Coarsening (gas diffusion from smaller to larger bubbles)

- ...a most unusual form of condensed matter
 - Like a gas:
 - volume \sim temperature / pressure
 - Like a liquid:
 - Flow without breaking
 - Fill any shape vessel
 - Under large force, bubbles rearrange their packing configuration
 - Like a solid:
 - Support small shear forces elastically
 - Under small force, bubbles distort but don't rearrange

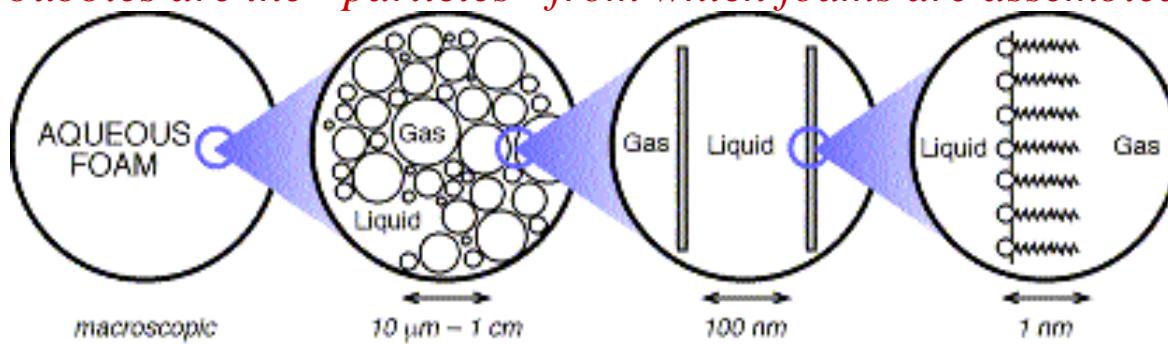


- **Everyday life:**
 - detergents
 - foods (ice cream, meringue, beer, cappuccino, ...)
 - cosmetics (shampoo, mousse, shaving cream, tooth paste, ...)
 - **Unique applications:**
 - firefighting
 - isolating toxic materials
 - physical and chemical separations
 - oil recovery
 - cellular solids
 - **Undesirable occurrences:**
 - mechanical agitation of multicomponent liquid
 - pulp and paper industry
 - paint and coating industry
 - textile industry
 - leather industry
 - adhesives industry
 - polymer industry
 - food processing (sugar, yeast, potatoes)
 - metal treatment
 - waste water treatment
 - polluted natural waters
 - *need to control stability and mechanics*
 - *must first understand microscopic structure and dynamics...*
- familiar!
 - important!

Condensed-matter challenge

- To understand the stability and mechanics of bulk foams in terms of the behavior at microscopic scales

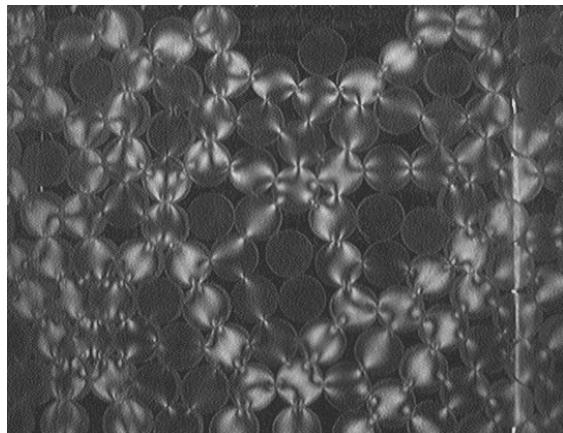
bubbles are the “particles” from which foams are assembled



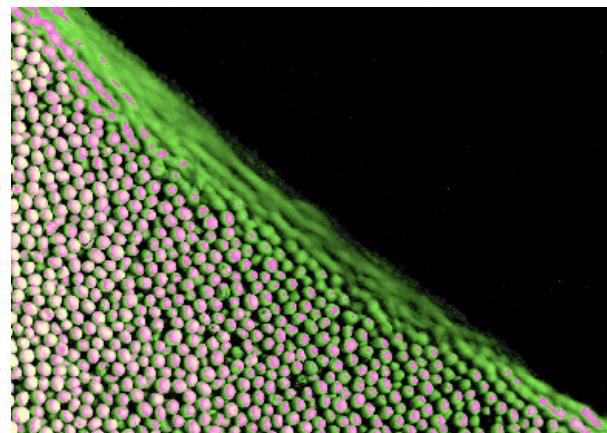
- Easy to relate surfactant-film and film-bubble behaviors
- Hard to relate bubble-macro behavior
 - Opaque: *no simple way to image structure*
 - Disordered: *no periodicity*
 - $k_B T \ll$ interaction energy: *no stat-mech.*
 - Flow beyond threshold: *no linear response*

- hard problems!
- new physics!

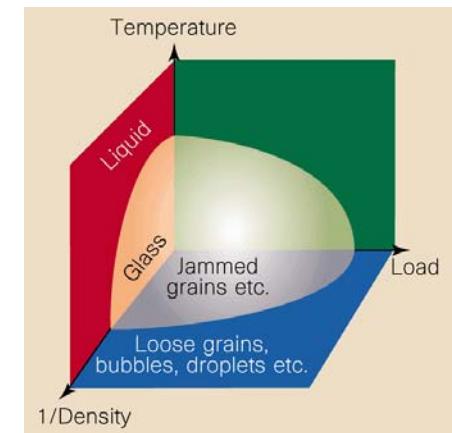
- Similar challenge for seemingly unrelated systems
 - Tightly packed collections of bubbles, droplets, grains, cells, colloids, fuzzy molecules, tectonic plates,....
 - jammed/solid-like: *small-force / low-temperature / high-density*
 - fluid/liquid-like: *large-force / high-temperature / low-density*



force-chains (S. Franklin)



avalanches (S.R. Nagel)



universality?



Foam Physics Today

- visit the websites of these Summer 2002 conferences to see examples of current research on aqueous foams
 - Gordon Research Conference on Complex Fluids
 - Oxford, UK
 - EuroFoam 2002
 - Manchester, UK
 - Foams and Minimal Surfaces
 - Isaac Newton Institute for Mathematical Sciences
 - Geometry and Mechanics of Structured Materials
 - Max Planck Institute for the Physics of Complex Systems



after these lectures, you should be in a good position to understand the issues being addressed & progress being made!



General references

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11. R. K. Prud'homme and S. A. Khan, ed., *Foams: Theory, Measurement, and Application*. Surfactant Science Series **57**, (Marcel Dekker, NY, 1996).
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15. L.J. Gibson and M.F. Ashby, *Cellular Solids: Structure and Properties* (Cambridge University Press, Cambridge, 1997).
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17. D. Weaire and S. Hutzler, *The Physics of Foams* (Clarendon Press, Oxford, 1999).
18. S.A. Koehler, S. Hilgenfeldt, and H.A. Stone, "A generalized view of foam drainage," *Langmuir* **16**, 6327-6341 (2000).
19. A.J. Liu and S.R. Nagel, eds., *Jamming and Rheology* (Taylor and Francis, New York, 2001).
20. J. Banhart and D. Weaire, "On the road again: Metal foams find favor," *Physics Today* **55**, 37-42 (July 2002).

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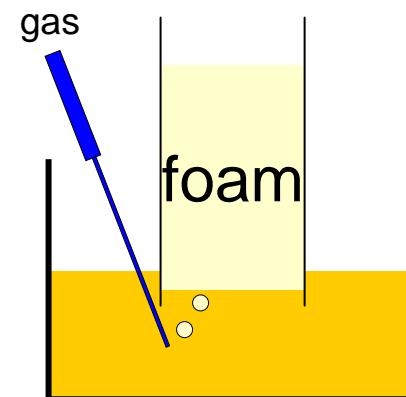
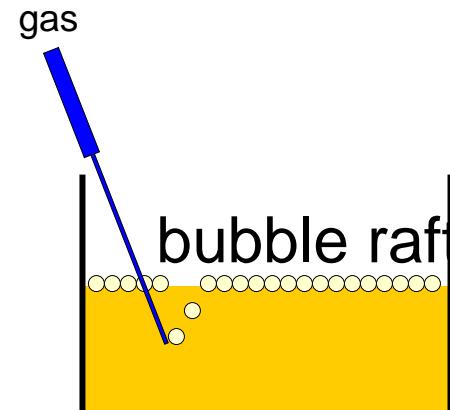
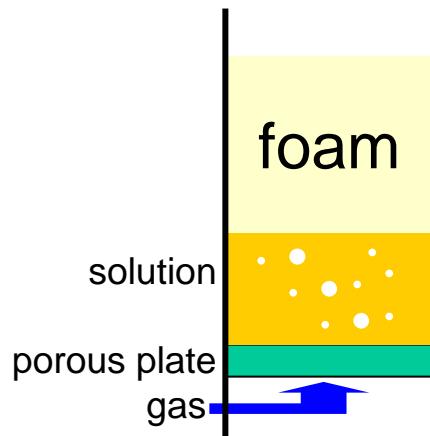
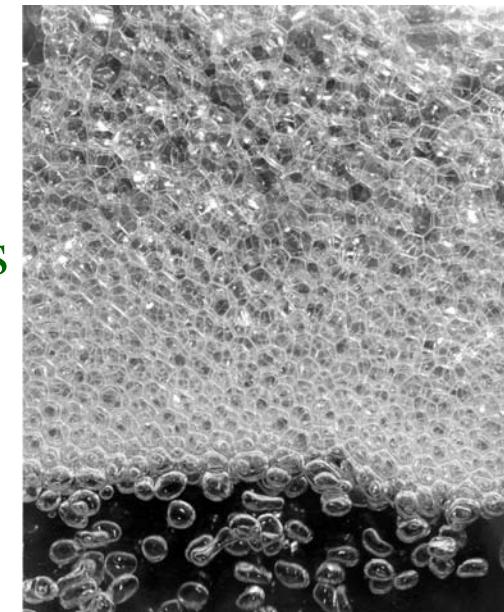
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Foam production I.

- Shake, blend, stir, agitate, etc.
 - Uncontrolled / irreproducible
 - Unwanted foaming of multicomponent liquids
- Sparge = blow bubbles
 - Polydisperse or monodisperse
 - Uncontrolled/non-uniform liquid fraction



Foam production II.

- *in-situ* release / production of gas

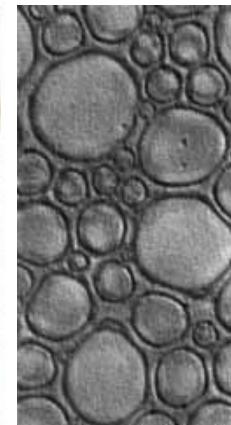
- nucleation

- eg CO₂ in beer



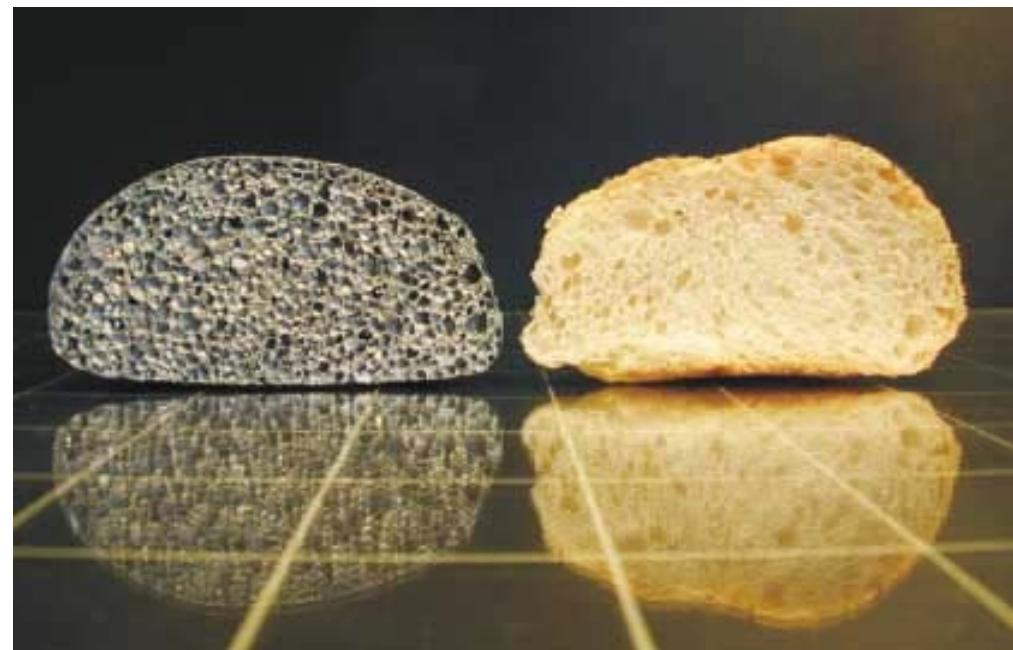
- aerosol:

- eg propane in shaving cream
 - small bubbles!



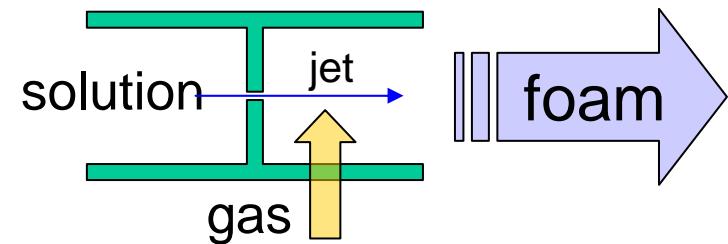
- active:

- eg H₂ in molten zinc
 - eg CO₂ from yeast in bread



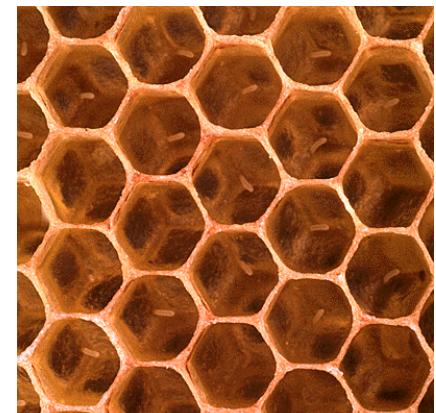
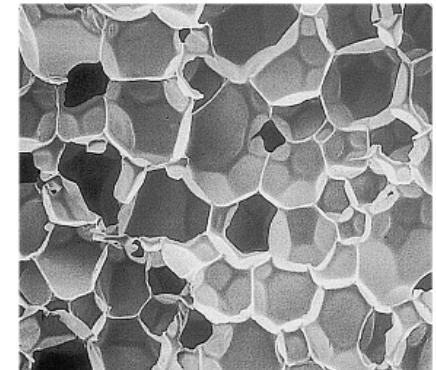
Foam production III.

- turbulent mixing of thin liquid jet with gas
 - vast quantities
 - small polydisperse bubbles
 - controlled liquid fraction
 - lab samples
 - firefighting
 - distributing pesticides/dyes/etc.
 - covering landfills
 - suppressing dust
 - ...



Foam production IV.

- many materials can be similarly foamed
 - nonaqueous liquids (oil, ferrofluids,...)
 - polymers (styrofoam, polyurethane,...)
 - metals
 - glass
 - concrete
- variants found in nature
 - cork
 - bone
 - sponge
 - honeycomb



Foams produced by animals

- spittle bug:



- cuckoo spit / froghoppers:



- stickleback-fish's nest

Foam production V.

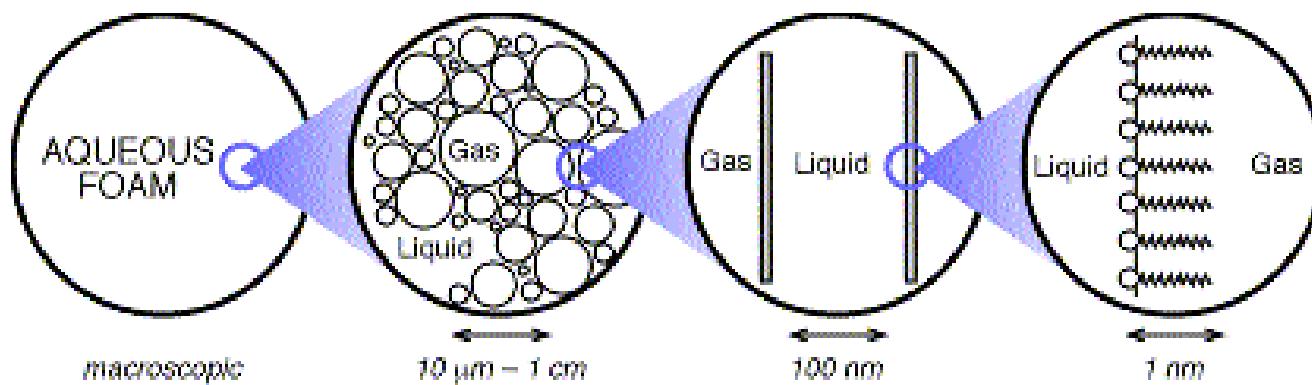
- antifoaming agents
 - prevent foaming or break an existing foam



- mysterious combination of surfactants, oils, particles,...

Microscopic behavior

- look at progressively larger length scales...



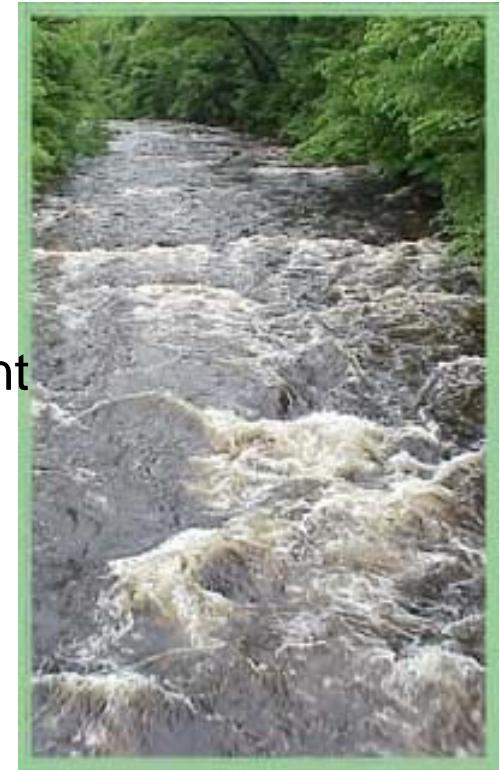
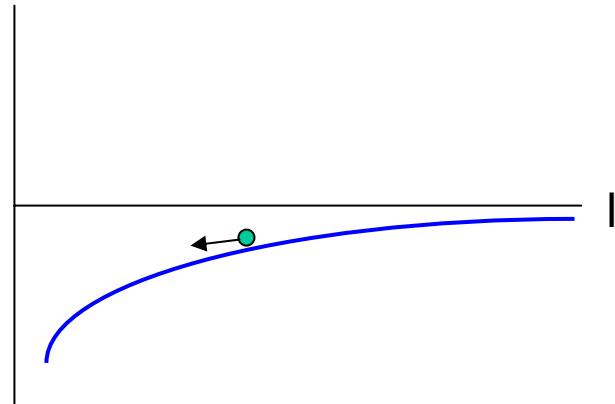
- surfactant solutions
- soap films
- local equilibrium & topology

- bubbles quickly coalesce – no foam
 - van der Waals force prefers monotonic dielectric profile; therefore, bubbles attract:



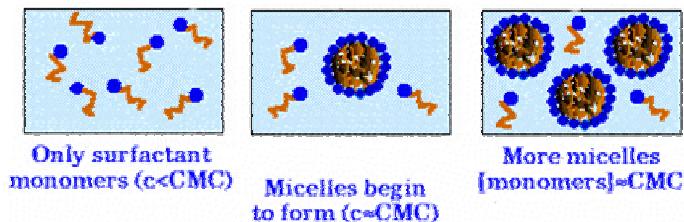
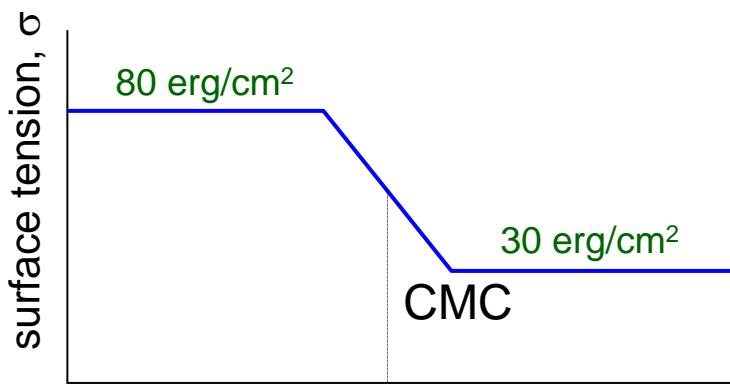
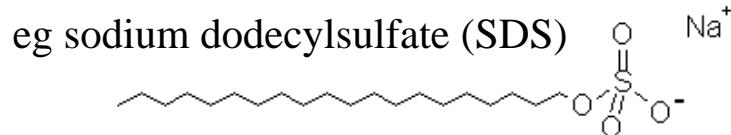
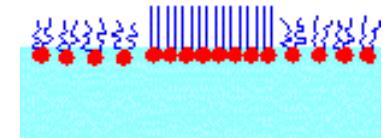
“effective interface potential”
is free energy cost per unit area:

$$V_{vdw}(l) = -A/12\pi l^2, \text{ A=Hamaker constant}$$



Surfactant solution

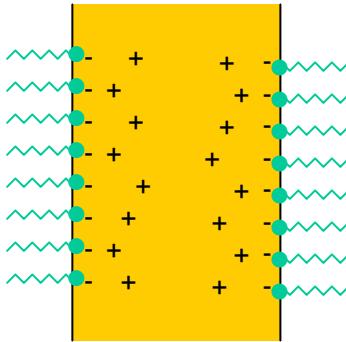
- **surface active** agent – adsorbs at air/water interface
 - head: hydrophilic (eg salt)
 - tail: hydrophobic (eg hydrocarbon chain)
 - look for good foams...
 - chain length: short enough that the surfactant is soluble
 - concentration: just above the “critical micelle concentration”



{NB: lower σ doesn't stabilize the foam..}

Electrostatic “double-layer”

- adsorbed surfactants dissociate, cause repulsion necessary to overcome van der Waals and hence stabilize the foam
 - electrostatic
 - entropic (dominant!)



free energy cost per unit area:

$$V_{DL}(l) = (64k_B T \rho / K_D) \text{Exp}[-K_D l],$$

ρ = electrolyte concentration

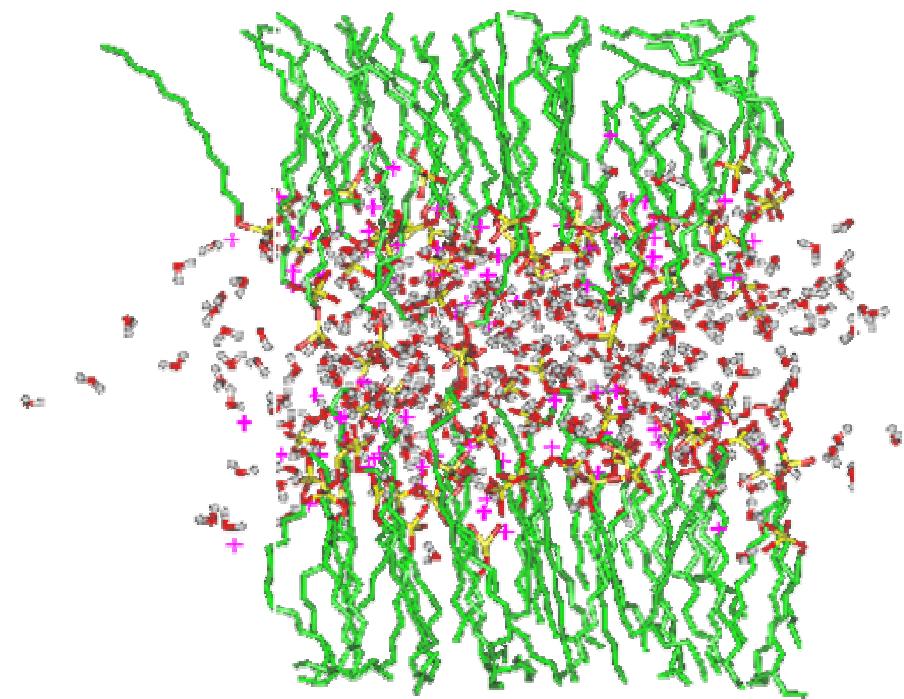
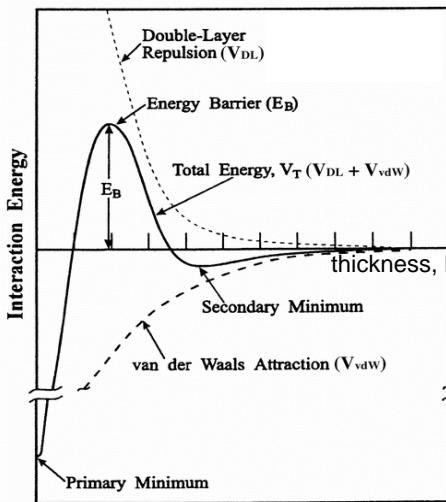
$K_D^{-1} \sim \rho^{-1/2}$ = Debye screening length

- NB: This is similar to the electrostatic stabilization of colloids

Soap film tension

- film tension / interface potential / free energy per area:

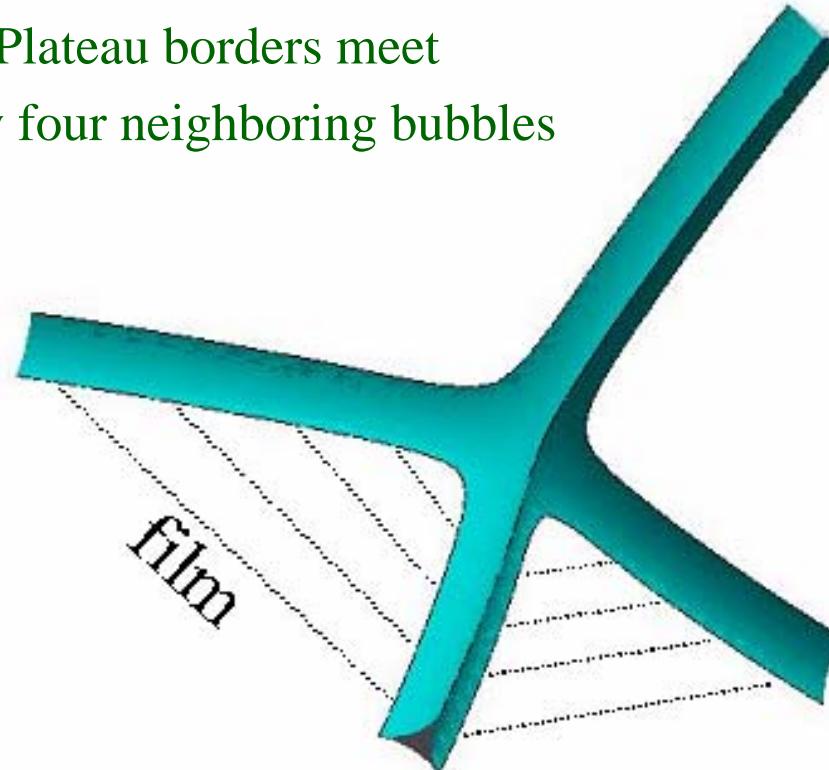
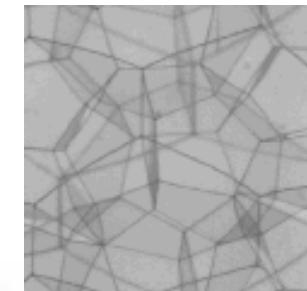
$$\gamma(l) = 2\sigma + V_{VDW}(l) + V_{DL}(l) \sim 2\sigma$$



- disjoining pressure: $\Pi(l) = -d\gamma/dl$
 - vanishes at equilibrium thickness, $l_{eq} \sim K_D^{-1}$ (30-3000 Å)

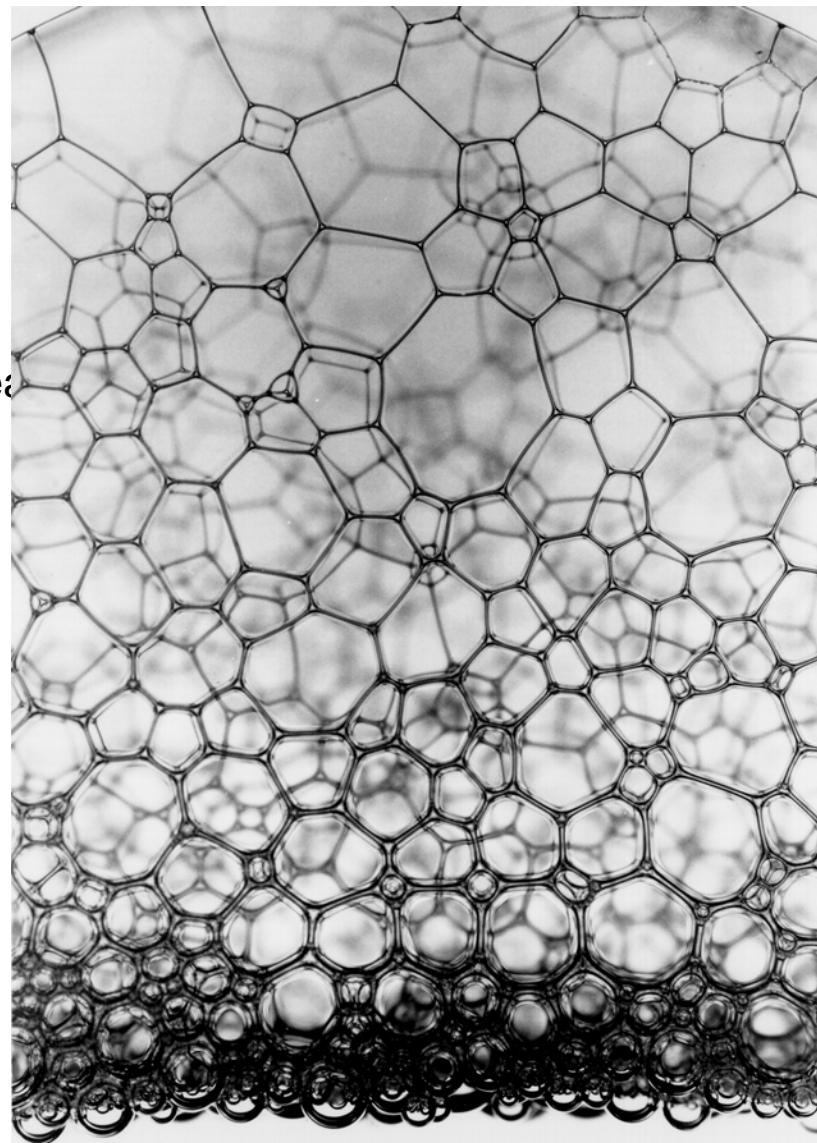
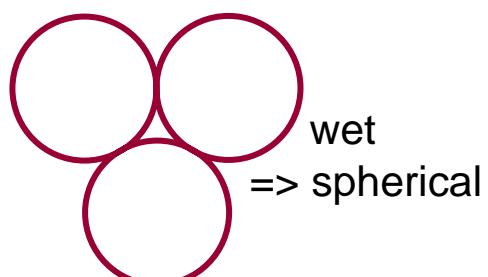
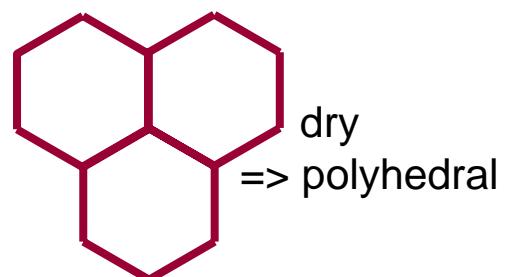
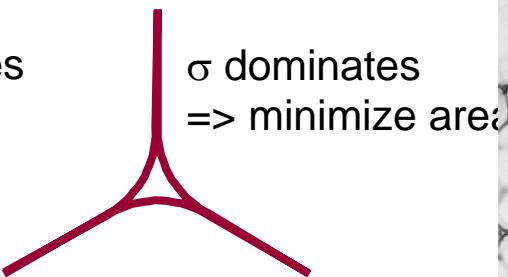
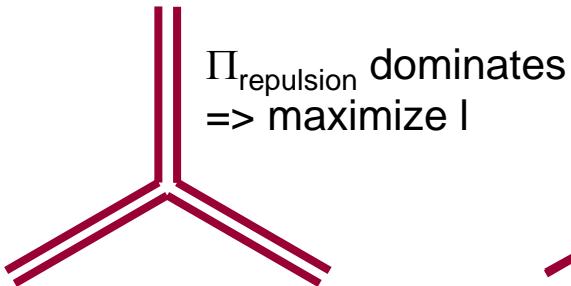
Film junctions

- Plateau border
 - scalloped-triangular channel where three films meet
 - the edge shared by three neighboring bubbles
- Vertex
 - region where four Plateau borders meet
 - the point shared by four neighboring bubbles



Liquid distribution

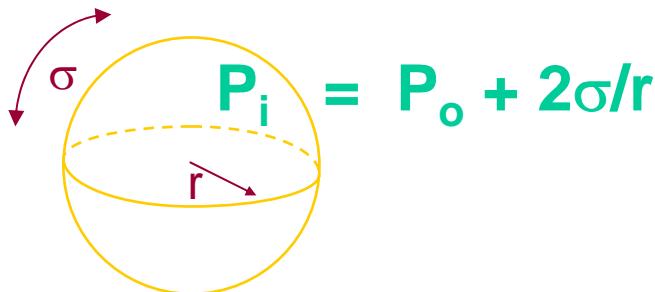
- division of liquid between films-borders-vertices
 - repulsion vs surface tension
 - wet vs dry



Laplace's law

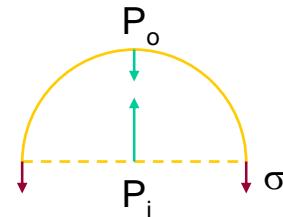
- the pressure is greater on the inside a curved interface

due to surface tension, $\sigma = \text{energy} / \text{area} = \text{force} / \text{length}$



$$P_i = P_o + 2\sigma/r$$

{in general, $\Delta P = \sigma(1/r_1 + 1/r_2)$ }



- forces on half-sphere:

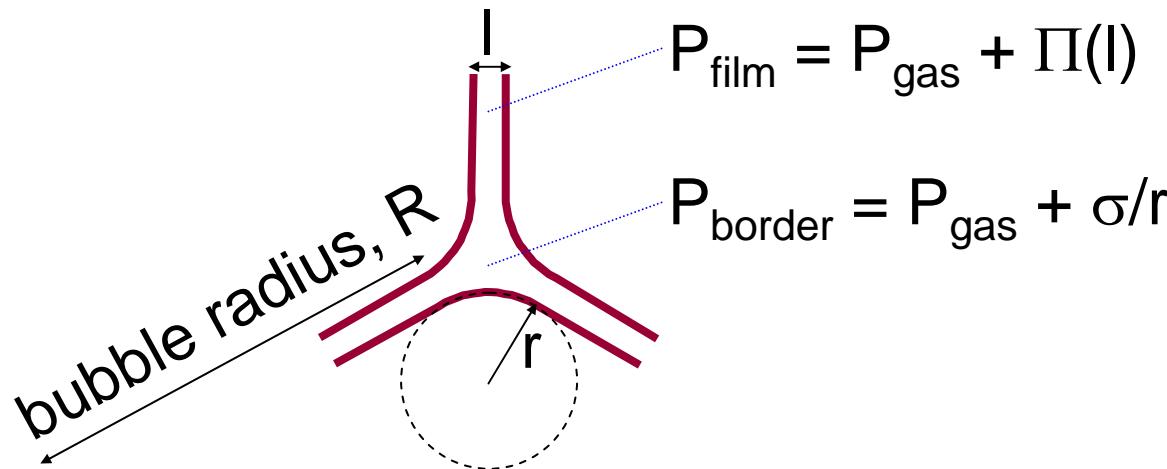
- $\sum F_{up} = P_i \pi r^2 - P_o \pi r^2 - 2\pi\sigma r = 0$

- energy change = pressure \times volume change:

- $dU = (\Delta P)4\pi r^2 dr$, where $U(r) = 4\pi r^2 \sigma$

Liquid volume fraction

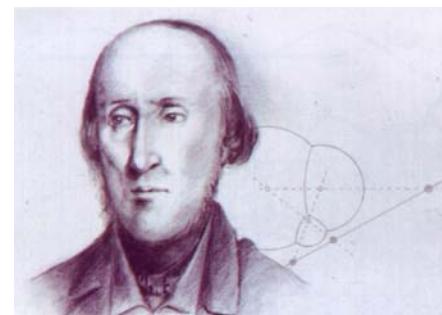
- liquid redistributes until liquid pressure is same everywhere



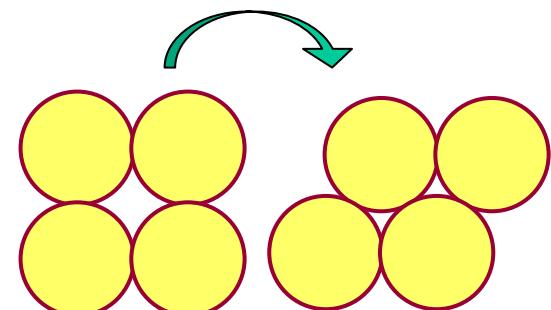
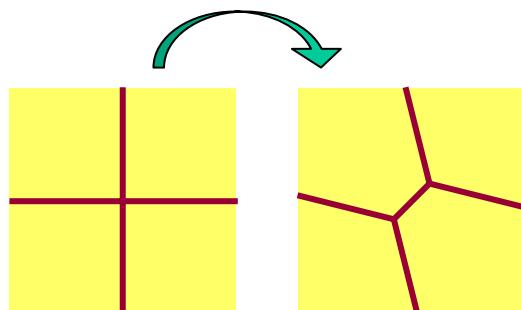
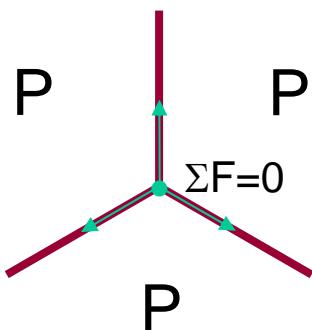
- typically: film thickness $l \ll$ border radius $r \ll$ bubble radius R
 - liquid volume fraction scales as $\varepsilon \sim (lR^2 + r^2R + r^3)/R^3 \sim (r/R)^2$
 - most of the liquid resides in the Plateau borders
 - PB's scatter light...
 - PB's provide channel for drainage...

Plateau's rules for dry foams

- for mechanical equilibrium:
 - i.e. for zero net force on a Plateau border,
 - zero net force on a vertex,
 - and $\sum \Delta P = 0$ going around a closed loop:

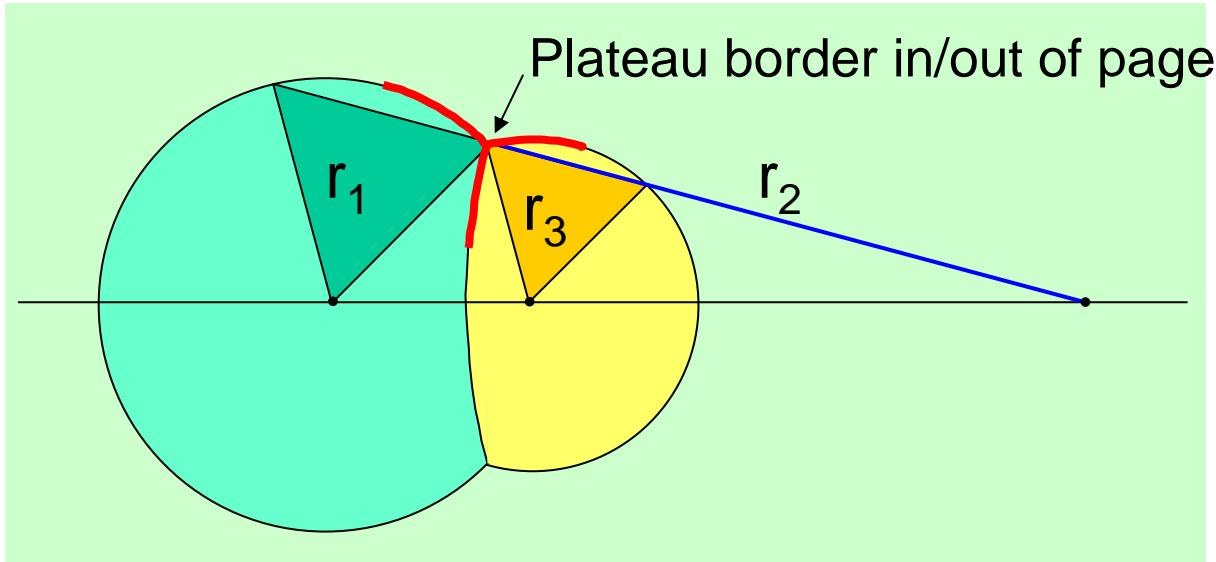


- (1) films have constant curvature & intersect three at a time at 120°
- (2) borders intersect four at a time at $\cos^{-1}(1/3)=109.47^\circ$
 - rule #2 follows from rule #1
 - both are obviously correct if the films and borders are straight:



Rule #1 for straight borders

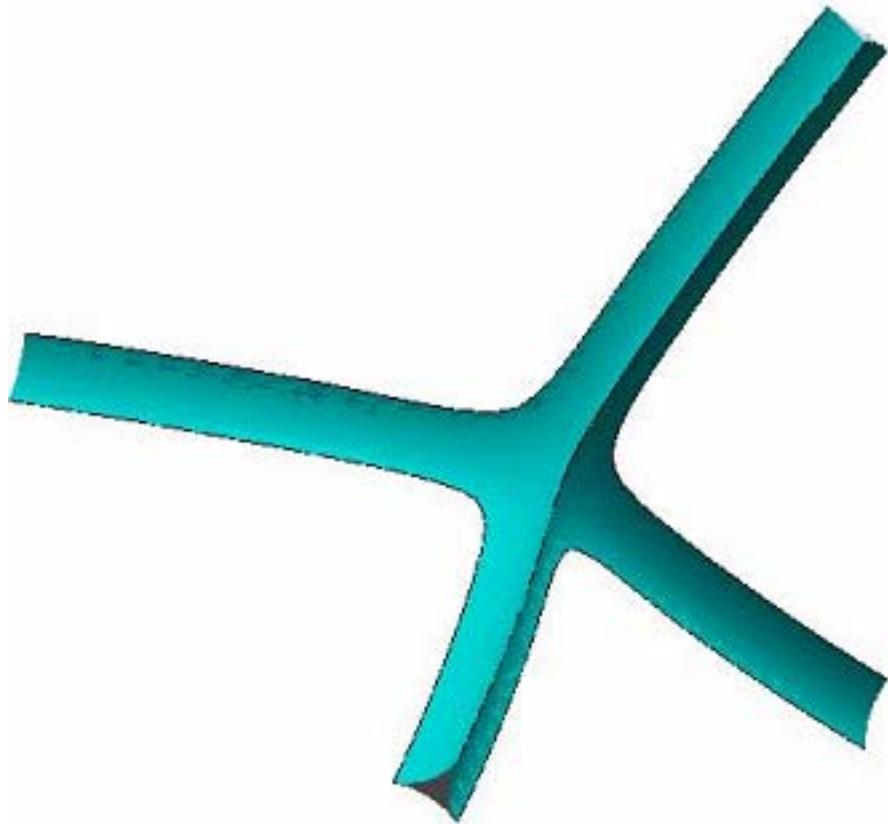
- choose r_1 and orientation of equilateral triangle
- construct r_2 from extension down to axis
- construct r_3 from inscribed equilateral triangle
 - NB: centers are on a line



- films meet at 120° (triangles meet at $60^\circ-60^\circ-60^\circ$, and are normal to PB's)
- similar triangles give $(r_1+r_2)/r_1 = r_2/r_3$, i.e. $1/r_1 + 1/r_2 = 1/r_3$ and so $\Sigma P=0$

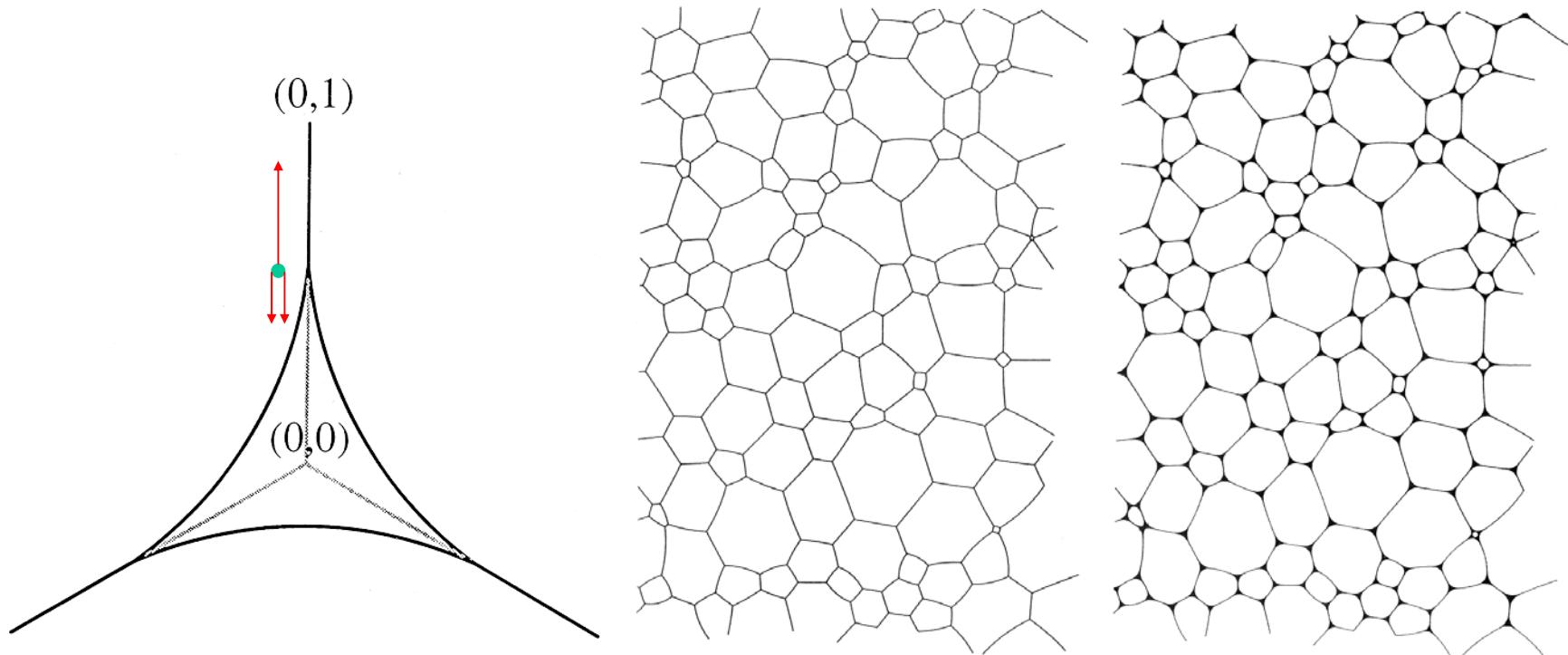
Curved Plateau borders

- proof of Plateau's rules is not obvious!
 - established in 1976 by Jean Taylor



Decoration theorem for wet foams

- for $d=2$ dimensions, an equilibrium wet foam can be constructed by *decorating* an equilibrium dry foam
 - can you construct an elementary proof?
 - PB's are circular arcs that join tangentially to film
 - theorem fails in $d=3$ due to PB curvature



- periodic foam structures
- disordered foam structures
 - experiment
 - simulation